



23rd ABCM International Congress of Mechanical Engineering December 6-11, 2015, Rio de Janeiro, RJ, Brazil

AVERAGE OF ELASTIC PROPERTIES OF THE CORNEA

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Abstract. Sometimes the analysis of the stress state, crack initiation and some others phenomenon in heterogeneous media is complicated. In order to facilitate the ulterior study of such problems, we apply the asymptotic homogenization method (AHM) to a general laminated shell composite and an equivalent homogeneous elastic problem with effective properties is obtained. The physical and material properties of the bio-composites are studied using the two scale of AHM. The heterogeneous elastic problem is solved. As examples of the application of the AHM the cornea is considered as a laminated shell bio-composite. The cornea is composed with three cell layers: the outer epithelial cells, stromal cells and endothelial cells. Between these layers are extracellular structures called Bowman's and Descemet's membranes. The cornea has a variable thickness. The polar coordinates system is introduced to study the geometry of this media. A function that describes the variation of the thickness is used to generate the model of the cornea. The effective elastic tensor is calculated by AHM and FEM considering two models, one where the thickness is constants along the cornea and other considering thickness variation. These two models by AHM and FEM are compared and the influence of the variation of the thickness.

Keywords: Effective elastic properties, Asymptotic Homogenization Method, Bio-composites, Laminated structures, Numerical analysis, Finite Element Methods

1. INTRODUCTION

The cornea possesses a unique combination of mechanical stiffness, strength, and optical transparency that enables it to serve as both a protective covering and the primary refractive component of the eye. These properties depend on the composition and orientation of the cornea, and differ from the central cornea to the limbus, (Nguyen and Boyce, 2011). The stress and the strain are related with the deformation and the geometry of the cornea, hence the importance of the study of the elastic properties of the cornea (Holmes *et al.*, 2001), (Garcia-Porta *et al.*, 2014), (Voorhies, 2003). Otherwise, the glaucoma and myopia are the most common problems that the cornea has on these days (Gatinel *et al.*, 2001), (Broman *et al.*, 2007), (González-Méijope *et al.*, 2008). The transplants and surgery have to be very precise proceeding, then knowledge of the effective properties become an important tool because in a certain kind of eye surgery, the human eyeball is deformed sustainably by the application of an elastic band. To many author talk about the geometry of the cornea and

its influence on the mechanical properties, (Donohue *et al.*, 1995), (Pinsky *et al.*, 2005), (Sigat *et al.*, 2004). Using the two scale Asymptotic Homogenization Method and Finite Elements Method, we determinate the effective properties of the cornea and compare the influence of its geometry on the results.

2. LINEAR ELASTIC PROBLEM IN THE CORNEA

In (Skacel and Bursa, 2014), (Genest, 2010), (Li and Tighe, 2006), (Pandolfi and Holzapfel, 2008), (Cardoso and Cowin, 2012), the stress-strain relationship in the cornea has an exponential behavior and the analysis of the elastic properties is made considering a non-linear model. Otherwise, beyond a certain stress, the $\sigma - \epsilon$ relationship becomes almost linear (Elsheikh *et al.*, 2011).

In order to reduce the analysis of the elastic theory problem in the cornea, a linear constitutive equation for the stress and the strain is considered, following the models presented in (Asher *et al.*, 2014), (Cabrera-Fernádez *et al.*, 2005), (Donohue *et al.*, 1995), (Fisher, 1971), (Anderson *et al.*, 2004).

The cornea is considered a solid periodic structure Ω bound by the surfaces Σ_1 and Σ_2 . Fixed a coordinate system $\mathbf{x} = (x_1, x_2, x_3)$, the geometry of the cornea is described by a function $\boldsymbol{\varrho}(\mathbf{x})$. For small deformations, the operator relating the strain $\boldsymbol{\epsilon}$ and the stress $\boldsymbol{\sigma}$ in the cornea is the Hooke's law,

$$\boldsymbol{\sigma} = \boldsymbol{C}(\mathbf{x}, \mathbf{y}) : \boldsymbol{\epsilon},\tag{1}$$

where C is the elastic tensor, regular in x and periodic in y; here $y = \rho(x)/\varepsilon$ is the fast or local variable and ε is a very small parameter which characterizes the periodicity. Following the Einstein's convention of summing over repeated indices, the elastic problem takes the form

$$\sigma^{ij}\|_j + f_i = 0 \text{ on } \Omega, \tag{2}$$

with boundary conditions

$$u_j = u_j^0 \text{ on } \Sigma_1, \qquad \sigma^{ij} \cdot n_j = S^i \text{ on } \Sigma_2, \tag{3}$$

where $\{\cdot\}\|_j$ denotes the covariant derivative with respect to the variable x_j , f_i are the components of the body forces, u_j are the components of the displacement, n_j is the normal vector in Σ_2 and the functions u_i^0 and S^i are the values of the displacement and the stress in Σ_1 and Σ_2 , respectively.

The general expression of the Cauchy's formula is

$$\epsilon_{mn} = \frac{1}{2} (u_m \|_n - u_n \|_m). \tag{4}$$

Taking into account the symmetry of the elastic tensor C^{ijkl} and substituting the equation (4) in equation (1), the contravariant components of the stress tensor are

$$\sigma^{ij} = C^{ijkl} (u_{k,l} - \Gamma^p_{kl} u_p), \tag{5}$$

where Γ_{kl}^{p} denotes the Christoffel's symbols of second kind.

Substituting equation (5) into equation (2), the general expression of the linear elastic problem can be obtained:

$$\left(\frac{\varrho_{k,j}}{\varepsilon}C^{ijmn}_{k} + C^{ijmn}_{,j} + \Gamma^{i}_{jk}C^{kjmn} + \Gamma^{j}_{jk}C^{ikmn}\right)(u_{m,n} - \Gamma^{r}_{mn}u_{r}) + C^{ijmn}\left(u_{m,nj} - \Gamma^{r}_{mn,j}u_{r} - \Gamma^{r}_{mn}u_{r,j}\right) + f^{i} = 0 \quad \text{on} \quad \Omega,$$
(6)

with the boundary conditions

$$u_i = u_i^0 \text{ on } \Sigma_1 \quad \left(C^{ijmn} \left(u_{m,n} - \Gamma_{mn}^r u_r \right) \right) n_j = S^i \text{ on } \Sigma_2, \tag{7}$$

where $\{\cdot\}_{j} = \frac{\partial}{\partial y_{j}} \{\cdot\}$ denotes the differentiation with respect to the fast or local curvilinear coordinate.

One of the main difficulties in solving problem (6)-(7) resides in the rapidly oscillating coefficients. In this paper we reduce it to an equivalent problem with homogenized coefficients, and thus obtain the effective properties of the cornea, by using the two-scale Asymptotic Homogenization Method (AHM) described next.

3. ASYMPTOTIC HOMOGENIZATION METHOD

The AHM is used to determine the homogenized problem and the effective properties of the solution of (6)-(7). The general expression for the two-scale asymptotic expansion is

$$u_m^{(\varepsilon)} = v_m + \varepsilon \left[N_{(0)m}^p v_p + N_{(1)m}^{lk} v_{l,k} \right] + o(\varepsilon), \tag{8}$$

where the local functions $N_m^{lk} \equiv N_m^{lk}(\mathbf{x}, \mathbf{y})$ and $N_{(0)m}^p \equiv N_{(0)m}^p(\mathbf{x}, \mathbf{y})$ are \mathbf{y} -periodic, and $\langle N \rangle = 0$, where $\langle \cdot \rangle$ is the average operator, defined by $\langle \cdot \rangle = \frac{1}{V_{\mathbf{Y}}} \int_{\mathbf{Y}} (\cdot) \sqrt{g} d\mathbf{y}$. Here \mathbf{Y} is the periodic cell, $V_{\mathbf{Y}}$ is its volume, and $g = \det([g_{ij}])$ where $[g_{ij}]$ is the metric tensor.

Substituting the expansion (8) into the heterogeneous problem (6)-(7) we consider the the coefficients for ε^{-1} , which must vanish:

$$\left(\varrho_{q,j}C^{ijlk} + \varrho_{p,n}C^{ijmn}N^{lk}_{(1)m|p}\varrho_{q,j}\right)_{|q} = 0,\tag{9}$$

$$\left(-\varrho_{l,j}\mathcal{C}^{ijmn}\Gamma^p_{mn} + \varrho_{l,j}\mathcal{C}^{ijmn}\varrho_{t,n}N^p_{(0)m|t}\right)_{|l} = 0.$$
(10)

Solving the differential equation (9), the values of the local functions $N_{(1)m|p}^{lk}$ can be determined. Equation (10) allows to calculate $N_{(0)m|t}^{p}$ as a function of $N_{(1)m|p}^{lk}$,

$$\left(\varrho_{l,j}\mathcal{C}^{ijmn}\varrho_{t,n}N^p_{(0)m|t}\right)_{|l} = -\Gamma^p_{mn}\left(\varrho_{q,s}\mathcal{C}^{ijks}N^{mn}_{(1)k|q}\varrho_{l,j}\right)_{|l}.$$
(11)

Applying the average operator to the coefficient of ε^0 and grouping the coefficients of $v_{k,jl}$, $v_{k,l}$ and v_l , we get

$$C_e^{ijkl} \equiv \hat{h}^{ijkl} = \left\langle C^{ijkl} + C^{ijmn} \varrho_{p,n} N^{kl}_{(1)m|p} \right\rangle, \tag{12}$$

as well as expressions for the homogenized coefficients,

$$\hat{h}^{ilk} = \hat{h}^{ijlk}_{,j} - \Gamma^l_{mn} \hat{h}^{ikmn} + \Gamma^i_{jn} \hat{h}^{njlk} + \Gamma^j_{jn} \hat{h}^{inlk},$$
(13)

$$\hat{h}^{il} = -\left(\Gamma^l_{mn}\hat{h}^{ijmn}\right)_{,j} - \left(\Gamma^i_{jk}\hat{h}^{kjmn} + \Gamma^j_{jk}\hat{h}^{ikmn}\right)\Gamma^l_{mn},\tag{14}$$

$$\widehat{h}_*^{ijl} = -\Gamma_{mn}^l \widehat{h}^{ijmn}.$$
(15)

Using these expressions we write the homogenized problem, consequence of (6)-(7), but now with slowly oscillating coefficients:

$$\hat{h}^{il}v_l + \hat{h}^{ikl}v_{l,k} + \hat{h}^{imln}v_{l,nm} + f^i = 0 \quad \text{on} \quad \Omega,$$
(16)

and with boundary conditions

$$v_i = u_i^0 \text{ on } \Sigma_1 \quad \left(\widehat{h}_*^{ijl} v_l + \widehat{h}^{ijlk} v_{l,k}\right) n_j = S^i \text{ on } \Sigma_2.$$

$$(17)$$

Now, the effective linear stress has the expression

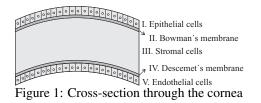
$$\sigma_e^{ij} = C_e^{ijkl} v_{k,l} + C_e^{ijk} v_k, \tag{18}$$

where C_e^{ijkl} are the effective coefficients (12), v_k is the solution of the homogenized problem (16)-(17) and C_e^{ijk} is

$$C_e^{ijk} = \left\langle C^{ijmn} N_{(0)m|n}^k - \Gamma_{mn}^k C^{ijmn} \right\rangle.$$
⁽¹⁹⁾

4. COMPOSITION AND GEOMETRY OF THE CORNEA

The cornea is laminate shell composite, with three cell layers: the outer epithelial cells (I), stromal cells (III) and endothelial cells (V). Between these layers are extracellular structures called Bowman's (II) and Descemet's (IV) membranes, (Anderson *et al.*, 2004; Dupps and Wilson, 2006). The layers (I)-(V) of the cornea have central thickness of $14\mu m$, $12\mu m$, $450\mu m$, $10\mu m$ and $14\mu m$, respectively. A cross-section of the cornea is shown in Fig. 1.



The spherical coordinate system (r, φ, θ) is used to describe the three-dimensional model of the cornea (Cardoso and Cowin, 2012), shown in Fig. 2a. In order to determine the effective properties of the cornea, we consider it rotationally invariant around its symmetry axis. In other words, its elastic properties have no variation in the θ -direction. As a consequence, we can use a two-dimensional model of the cornea, which in turn can be described by means of the polar coordinate system (r, φ) .

In our calculations we have used the following values of the cornea parameters: its thickness is approximately $T_1 = 0.5mm$ at its apex and $T_2 = 0.67mm$ at the limbus, (Pandolfi and Holzapfel, 2008; Ruberti *et al.*, 2011); the length of the half cross-section of the cornea is $D_1 = 6mm$, with $D_2 = 2.5mm$ of height. The inner radio in the center of the cornea is $R_1 = 7.56mm$, at the limbus the radio is $R_2 = 7.51mm$ and the angle between the radios is $\alpha = 0.8235$ (see Fig. 2b), (Cabrera-Fernádez *et al.*, 2005).

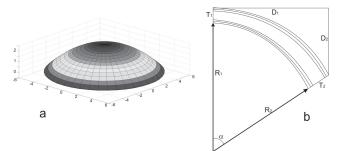


Figure 2: Geometry of the cornea: a) Three dimensional representation of the cornea given by the function (20). b) Geometry of the cross section of the cornea

Taking into account the values of the parameters R_1 , R_2 , D_1 , D_2 , T_1 , T_2 and α , we can use the following function that describes the geometry of the cornea:

$$\varrho(r,\varphi) = r + (m \cdot r + n) \left(\varphi - \pi/2\right)^2,\tag{20}$$

with m = 0.5014, n = -3.8643, $r \in [7.56, 8.06]$ and $\varphi \in [\pi/2, \pi/2 + \alpha]$.

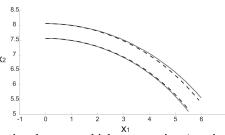


Figure 3: Comparison of the two-dimensional cornea: thickness varying (continuous line) and constant thickness along the cornea structure (discontinuous line), given by (20) and (21), respectively.

Alternatively, we consider also the constant-thickness model, corresponding to

$$\varrho(r,\varphi) = r. \tag{21}$$

Using the polar coordinates system both two-dimensional models of the cornea can be compared, see Fig. 3. The positions are given by the equations

$$x_1 = \rho \cos(\varphi), \quad x_2 = \rho \sin(\varphi),$$
(22)

where ρ is as in (20) and (21).

5. EFFECTIVE ELASTIC PROPERTIES OF THE CORNEA

The layers which compose the cornea are isotropic materials; the cornea is oriented un such a way that the fast variation of its properties corresponds to the y-direction, where $y = \rho/\varepsilon$ and ρ is given either by (20) or (21). In the polar coordinate system the metric tensor is $g_{ij} = \text{diag}[1, r^2, 1]$ and the general expression for the isotropic components of C^{ijkl} is

$$C^{ijkl} = \lambda g^{ij} g^{kl} + \mu (g^{ik} g^{jl} + g^{il} g^{jk}),$$
(23)

where $[g^{ij}] = [g_{ij}]^{-1}$.

Considering the polar metric tensor, the nonzero Christoffel's symbols are

$$\Gamma_{22}^1 = -r \quad \Gamma_{12}^2 = \Gamma_{21}^2 = \frac{1}{r}.$$
(24)

The elastic properties for the different layers of the cornea (see Fig. 1) are given in Table 1. These values were taken from (Cabrera-Fernádez et al., 2005). The elastic tensor (23) can be determined using the input parameters of Table 1.

Table 1: Mechanical properties of the cornea.				
Layers	Young's Modulus	Poisson's ratio		
Epithelial cells (I)	0.622	0.49		
Bowman's membrane (II)	0.275	0.49		
Stromal cells (III)	0.279	0.49		
Descemet's membrane (IV)	0.304	0.49		
Endothelial cells (V)	0.331	0.49		

Considering $y = \rho/\varepsilon$ one-dimensional, the local problems (9) and (10) take the following expressions, respectively:

$$\frac{\partial}{\partial y} \left(\varrho_{,j} C^{ijlk} + \varrho_{,n} C^{ijmn} \frac{\partial N^{lk}_{(1)m}}{\partial y} \varrho_{,j} \right) = 0, \tag{25}$$

$$\frac{\partial}{\partial y} \left(-\varrho_{,j} \mathcal{C}^{ijmn} \Gamma^p_{mn} + \varrho_{,j} \mathcal{C}^{ijmn} \varrho_{,n} \frac{\partial N^p_{(0)m}}{\partial y} \right) = 0.$$
(26)

In order to obtain the effective coefficients, the value of the local function $\partial N_{(1)m}^{lk}/\partial y$ is determined by integrating the local problem (25),

$$\frac{\partial N_{(1)m}^{lk}}{\partial y} = \left(A^{ilk} - \varrho_{,p}C^{iplk}\right) \left(\varrho_{,n}C^{ijmn}\varrho_{,j}\right)^{-1},\tag{27}$$

where the tensor $A^{ilk} \equiv A^{ilk}(\mathbf{x})$. With account of the periodic condition of the local function N_m^{lk} , the average $\left< \partial N_{(1)m}^{lk} / \partial y \right> = 0$, and the function A^{ilk} yields

$$A^{ilk} = \left\langle \varrho_{,p} C^{iplk} \left(\varrho_{,n} C^{ijmn} \varrho_{,j} \right)^{-1} \right\rangle \left\langle \left(\varrho_{,n} C^{ijmn} \varrho_{,j} \right)^{-1} \right\rangle^{-1}.$$
⁽²⁸⁾

Finally, functions $\partial N_{(1)m}^{lk}/\partial y$ are obtained by using the values of ρ and C^{ijkl} , given in (20) and (23), respectively, in the equations (27) and (28).

In this paragraph, the numerical method proposed to solve problems (25), (26) is described. Since the problem is composed of two coupled PDE's with singular (discontinuous) rigidity tensor, a P_1 (piecewise linear) finite element method (FEM) is used on a weak formulation of (25), (26).

Due to the properties of the finite element discretization, after the computation of an approximation of the solution $N_{(0)m}^p$ and $N_{(1)m}^{lk}$ using the FEM method, the integrals (see formula (12) for example) for the calculus of the effective coefficients are replaced by sums over each element contribution.

In what follows we use the super-indices t_v and t_0 in order to denote the computed coefficients corresponding to the thickness functions (20) and (21), respectively.

In (Anderson et al., 2004), (Asher et al., 2014), (Pandolfi and Holzapfel, 2008), (Lanchares et al., 2008), (Li and Tighe, 2006) and (Ruberti et al., 2011) the variable thickness along animal or human lens is considered. Otherwise, (Broman et al., 2007), (Last et al., 2012), (Donohue et al., 1995) and (Huang et al., 2002) a constant average thickness in the cornea's structure is analyzed. In this work, we address and compare both cases in order to analyze the influence of the thickness variation on the effective properties.

We also compare the values of the effective coefficients C_e^{ijkl} obtained by AHM and the finite element method (FEM) in the both models. The effective coefficients for t_0 , found by AHM and FEM, are shown in Table 2 for the values r = 7.56 and $\varphi = \pi/2$.

In the case of the varying thickness (denoted by the superscript t_v), the coefficients do depend on r and φ (see Tables

3 and 4). Thus, in Table 2 we present also the distance between $\hat{C}_{e}^{t_0}$ and $C_{e}^{t_v}$. We present also the values of the components C_{e}^{1111} (Table 3) and C_{e}^{1112} (Table 4) of the elastic tensor $C_{e}^{t_v}$ for the case of the varying thickness.

Coefficients	$m{C}_{e}^{t_{0}}$ (AHM)	$oldsymbol{C}_{e}^{t_{0}}$ (FEM)	$\ \boldsymbol{C}_{e}^{t_{0}}-\boldsymbol{C}_{e}^{t_{v}}\ $ (AHM)
C_e^{1111}	4.8812	4.8814	8×10^{-4}
C^{1122}	0.0821	0.0836	8×10^{-8}
C^{1133}	4.6898	4.6936	2×10^{-6}
C_{e}^{2222}	0.0015	0.0015	2×10^{-9}
C^{2233}	0.0821	0.0821	4×10^{-8}
$\begin{array}{c} C_e \\ \hline C_e^{3333} \\ e \end{array}$	4.8886	4.8888	3×10^{-15}
C_e^{1112} C^{2212}	0	0	1×10^{-8}
C_{e}^{2212}	0	0	3×10^{-7}
$\begin{array}{c} C_e \\ \hline C_e^{3312} \\ \hline \end{array}$	0	0	1×10^{-5}
C^{2323}	0.0017	0.0017	2×10^{-8}
C^{1323}	0	0	6×10^{-6}
C_{e}^{1313}	0.0957	0.0957	1×10^{-6}
C_{e}^{1212}	0.0017	0.0017	8×10^{-8}

Table 2: Effective coefficients $C_e^{t_0}$ calculated by AHM and FEM. Difference between $C_e^{t_0}$ and $C_e^{t_v}$.

Table 3: Components C_e^{1111} of the elastic tensor $\boldsymbol{C}_e^{t_v}$ for different values of r and φ .

$\varphi \rightarrow$	1.5708		1.9002		2.3943	
$r\downarrow$	AHM	FEM	AHM	FEM	AHM	FEM
7.56	4.88123	4.88133	4.88098	4.88132	4.88083	4.88129
7.76	4.88123	4.88133	4.88132	4.88134	4.88137	4.88133
8.06	4.88123	4.88133	4.88182	4.88127	4.88211	4.88221

Table 4: Components C_e^{1112} of the elastic tensor $C_e^{t_v}$ for different values of r and φ .

$\varphi \rightarrow$	1.57	708	1.9002		2.3943	
$r\downarrow$	AHM	FEM	AHM	FEM	AHM	FEM
7.56	0	0	0.0243×10^{-4}	0.0272×10^{-4}	0.0479×10^{-4}	0.0535×10^{-4}
7.76	0	0	-0.0088×10^{-4}	-0.0098×10^{-4}	-0.0172×10^{-4}	-0.0193×10^{-4}
8.06	0	0	-0.0584×10^{-4}	-0.0653×10^{-4}	-0.1148×10^{-4}	-0.1285×10^{-4}

We can appreciate a very good coincidence between the effective properties calculated by AHM and FEM. It should be taken into account that in the present situation even a small difference in the effective coefficients (for varying and constant thicknesses), as in the third column of Table 2, renders important difference in the anisotropy of the material. Namely, for the constant thickness with isotropic layers the homogenized media exhibits transversely isotropic behavior, but in the case with varying thickness the material has a monoclinic symmetry.

A important comparison between the calculation of the effective coefficients, obtained by AHM and FEM, is the time of the computational operations. The values of Table 2 have been reproduce using AHM in a average time of 0.051 seconds, however, the average time reported by FEM was 1.6 seconds.

6. SOLVING THE HOMOGENIZED PROBLEM OF THE CORNEA

An external uniform force is applied on a two-dimensional cornea, simulating the intra-ocular pressure (IOP). This force **f** occurs in the *r*-direction and has contravariant components ($\cos \varphi, 0, 0$) with respect to the basis of the tangent vectors \mathbf{g}_i . This force crates an internal displacement $\mathbf{u}|_{\Sigma_1} = (0, 0, 0)$ and the traction $\mathbf{S}|_{\Sigma_2}$ that has the same contravariants components of the force **f**.

The influence of the thickness variation in the stress tensor is analyzed. We consider a particular case of the elastic problem. The variation of the properties is in the x_1 direction and the only non zero displacement is u_1 . Under these assumptions, the homogenized problem (16)-(17) is

$$\hat{h}^{11}v_1 + \hat{h}^{111}v_{1,1} + \hat{h}^{1111}v_{1,11} + \cos\varphi = 0 \quad \text{on} \quad r \in (7.56, 8.06),$$
(29)

with the boundary conditions

$$v_1 = 0 \text{ on } r = 7.56 \quad \left(\hat{h}_*^{111}v_1 + \hat{h}^{1111}v_{1,1}\right)n_1 = \cos\varphi \text{ on } r = 8.06.$$
 (30)

for $\varphi \in [\pi/2, \pi/2 + \alpha]$.

The resulting coefficients \hat{h}^{1111} appear in Table 2; the corresponding coefficients \hat{h}^{111} , \hat{h}^{11} and \hat{h}^{111}_* are determined using equations (13)-(15).

The problem (29)-(30) is solved at different points of the cornea. The effective stress is calculated using the equations (18) and (19) for the angles $\varphi = \pi/2$ (corneal apex) $\varphi = 1.9002$ (a point between the apex and the limbus) and $\varphi = \pi/2 + \alpha$ (limbus of the cornea).

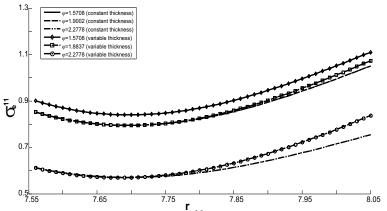


Figure 4: Comparison of the behavior of the effective stress σ_e^{11} at different points of the cornea with constant and variable thickness along the structure.

The component σ_e^{11} of the effective stress is varying along the cornea and has a non-linear behavior. At the center of the cornea ($\varphi = 1.5708$) the component σ_e^{11} for variable (\diamond) and constant thickness(–) coincide. However, the effective stress has different behavior for other values of φ , effect being more apparent at the limbus, ($\varphi = 2.2778$). Notice that the highest stress is located at the center of the cornea and the lowest one corresponds to the limbus.

7. CONSLUSIONS

The AHM is a powerful tool to determine the effective properties of the cornea quickly and with high precision. The methodology bring the possibility to aboard problems like the analysis of the bio-mechanics properties of the cornea after an instra-stromal ring segments implantation (Martínez-Finkelshtein *et al.*, 2009), (Martínez-Finkelshtein *et al.*, 2011), or the study of the variation of the corneal thickness in a healthy eye as function of the anterior face of the cornea and inner ocular pressure (IOP).

The difference of the length along the cornea is studied and its influence in the mechanical behavior of the cornea as a laminated shell bio-composite. The analytic expressions of the effective coefficients are very important result of this work. The numerical values are compared with the effective coefficients obtain by FEM, very good coincidence is reported and significant difference between the time of the computational operations.

ACKNOWLEDGEMENTS

The authors gratefully acknowledge to the project SHICHAN, supported by FSP (Cooperation Scientifique Franco-Cubaine) Cuba 2011-26, PROJET N° 29935XH and to the project Composite Materials from University of Havana. RRR thanks to Department of International Relations - DRI - Universidade Federal Fluminense - UFF, Edital Programa de Apoio a Projetos de Internacionalização da UFF - PIUFF - Edição 2014 for the financial suport during his stay at UFF as a visiting professor in 2015. AC gratefully thank the institutional support provided by INCT- MACC (National Institute of Science and Technology in Computer Science Medical Assistance and the Brazilian agency CNPq Project number 302298/2012-6. The fifth author (AMF) was partially supported by MICINN of Spain and by the European Regional Development Fund (ERDF) under grant MTM2011-28952-C02-01, by Junta de Andalucía (Excellence Grant P11-FQM-7276 and the research group FQM-229), and by Campus de Excelencia Internacional del Mar (CEIMAR) of the University of Almería. This work was completed during a visit of AMF to the Department of Mathematics of the Havana University, whose hospitality he wishes to acknowledge.

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